Momentum conserved GNN

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1 Introduction

A GNN takes as inputs initial node and edge features, $\mathbf{h}_{i}^{(0)}$ and $\mathbf{e}_{ij}^{(0)}$ respectively, and updates them after *L* layers, to get $\mathbf{h}_{i}^{(L)}$ and $\mathbf{e}_{ij}^{(L)}$. Then, the output acceleration \mathbf{a}_{i} is computed from them. We assume that external forces (e.g. gravity) are applied afterwards. Therefore, \mathbf{a}_{i} only includes the contribution from interaction with other particles. Naively, we can just perform for instance

$$\mathbf{a}_i = \phi(\mathbf{h}_i^{(L)}) \tag{1}$$

To ensure conservation of momentum, we need to ensure that

$$\sum_{i} \mathbf{a}_{i} = 0 \tag{2}$$

Note that Eq. 1 does not guarantees momentum conservation.

To fulfill that condition, we can write that as

$$\mathbf{a}_i = \sum_{j \in \mathcal{N}_i} \mathbf{f}_{ji} \tag{3}$$

where $\mathbf{f}_{ij} = -\mathbf{f}_{ji}$ is antisymmetric. The above could be understood as forces acting on node *i*. To show that the above prescription actually allows to conserve momentum, it is useful to define the function θ_{ji} such that it is 1 if $j \in \mathcal{N}_i$ and 0 otherwise. Note that *j* being in the neighborhood of *i* implies that *i* is also in the neighborhood of *j* (since the graph is bidirectional), so θ_{ij} is symmetric: $\theta_{ij} = \theta_{ji}$. Therefore,

$$\sum_{i} \mathbf{a}_{i} = \sum_{i} \sum_{j \in \mathcal{N}_{i}} \mathbf{f}_{ji} = \sum_{i,j} \theta_{ji} \mathbf{f}_{ji} = \frac{1}{2} \sum_{i,j} \left(\theta_{ij} \mathbf{f}_{ij} + \theta_{ji} \mathbf{f}_{ji} \right) = \frac{1}{2} \sum_{i,j} \mathbf{f}_{ij} \left(\theta_{ij} - \theta_{ji} \right) = 0$$
(4)

The feature vector \mathbf{f}_{ij} could be given by any antisymmetric combination of node or edge features. We propose the following form:

$$\mathbf{f}_{ij} = \mathbf{e}_{ij}^{(L)} - \mathbf{e}_{ji}^{(L)} \tag{5}$$

In these cases, $\mathbf{e}_{ij}^{(L)}$ should have the shape $(N_e, 3)$, with N_e the number of edges. What about the boundaries?

To account for external forces and boundary forces, we could add updated node features.

1.1 Energy conservation

$$E = \sum_{i} T_i + U_i \tag{6}$$

$$\frac{dE}{dt} = \frac{1}{m} \sum_{ij} \left(\mathbf{p}_i \cdot \mathbf{f}_{ji} - \frac{1}{2} \mathbf{f}_{ji} \cdot (\mathbf{p}_i - \mathbf{p}_j) \right) = 0 \tag{7}$$

Symmetries $\mathbf{2}$

• Euclidean

$$\mathbf{x}_i \rightarrow \mathbf{R}\mathbf{x}_i + \mathbf{T}$$
 (8)

$$\mathbf{v}_i \rightarrow \mathbf{R}\mathbf{v}_i$$
 (9)

$$\mathbf{a}_i \rightarrow \mathbf{R}\mathbf{a}_i$$
 (10)

 \bullet Galilean

$$\mathbf{x}_i \rightarrow \mathbf{R}\mathbf{x}_i + \mathbf{V}t$$
 (11)

$$\mathbf{v}_i \rightarrow \mathbf{R}\mathbf{v}_i + \mathbf{V}$$
 (12)

$$\mathbf{a}_i \rightarrow \mathbf{R}\mathbf{a}_i$$
 (13)

3 Integration

$$\mathbf{a}_i^t = GNN(\{\mathbf{x}_i^t, \mathbf{v}_i^t\}_{i \in G}) \tag{14}$$

• Symplectic Euler

$$\mathbf{v}_{i}^{t+1} = \mathbf{v}_{i}^{t} + \mathbf{a}_{i}^{t} \Delta t$$

$$\mathbf{x}_{i}^{t+1} = \mathbf{x}_{i}^{t} + \mathbf{v}_{i}^{t+1} \Delta t = \mathbf{x}_{i}^{t} + \mathbf{v}_{i}^{t} \Delta t + \mathbf{a}_{i}^{t} \Delta t^{2}$$
(15)
(15)
(15)

$$\mathbf{x}_{i}^{t+1} = \mathbf{x}_{i}^{t} + \mathbf{v}_{i}^{t+1}\Delta t = \mathbf{x}_{i}^{t} + \mathbf{v}_{i}^{t}\Delta t + \mathbf{a}_{i}^{t}\Delta t^{2}$$
(16)

• Leapfrog

$$\mathbf{v}_{i}^{t+1} = \mathbf{v}_{i}^{t} + \frac{1}{2} \left(\mathbf{a}_{i}^{t} + \mathbf{a}_{i}^{t+1} \right) \Delta t$$
(17)

$$\mathbf{x}_{i}^{t+1} = \mathbf{x}_{i}^{t} + \mathbf{v}_{i}^{t} \Delta t + \frac{1}{2} \mathbf{a}_{i}^{t} \Delta t^{2}$$
(18)

Note that in this case one computes first \mathbf{x}_i^{t+1} and the use that result to get \mathbf{a}_i^{t+1} and \mathbf{v}_i^{t+1} .

Symplectic Euler and leapfrog fulfill equivariance under Galilean and Euclidean transformations.